

The Photoelectric Effect - an Inelastic Scattering Process

Process



$$\hat{H}(t) = \hat{H}_0 + \hat{H}_1(t)$$

↑ neutral atom (e.g. Hydrogen)

$$\hat{H}_1(t) = -\left(\frac{e}{m}\right) \hat{\vec{A}} \cdot \hat{\vec{p}} + \left(\frac{e}{m}\right)^2 \hat{\vec{A}} \cdot \hat{\vec{A}}$$

↑ rad. ↑ atom

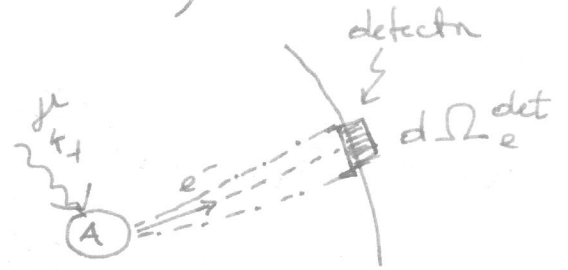
= 0 (single photon process)

↑ $\hat{\vec{A}}$ - ops. on photon (radiation states)

↑ $\hat{\vec{p}}$ - op. on electron (atomic states)

$$\hat{\vec{A}}(\vec{r}, t) = \sum_{\kappa', \lambda'} \left(\frac{\hbar}{2\epsilon_0 \omega_{\kappa'} V} \right)^{1/2} \left(\hat{a}_{\kappa', \lambda'} e^{i(\vec{k}' \cdot \vec{r} - \omega_{\kappa'} t)} + \hat{a}_{\kappa', \lambda'}^\dagger e^{-i(\vec{k}' \cdot \vec{r} - \omega_{\kappa'} t)} \right)$$

FGR for photoelectron:



$$d\Gamma_{fi} = \frac{2\pi}{\hbar} |V_{fi}|^2 \rho_e(E_e) d\Omega_e \int [E_{A^+} + E_e - (E_A + \hbar\omega_\kappa)] \quad (1)$$

> density of (final) states:

$$\rho_e(E_e) = \frac{V}{(2\pi\hbar)^3} \rho_e^2 \left(\frac{d\rho_e}{dE_e} \right) = \frac{V \rho_e^2}{(2\pi\hbar)^3} \frac{1}{\left(\frac{dE_e}{d\rho_e} \right)} = \frac{V \rho_e^2}{(2\pi\hbar)^3} \frac{1}{\frac{d}{d\rho} \left(\frac{\rho^2}{2m_e} \right)} = \frac{V \rho_e^2}{(2\pi\hbar)^3} \frac{m_e}{\rho_e}$$

$$= \frac{V}{(2\pi\hbar)^3} m_e \rho_e \quad (2)$$

> Transition matrix element:

$$V_{fi} \equiv \langle \phi_f | \hat{H}_1 | \phi_i \rangle$$

$$|\phi_i\rangle \equiv |A\rangle \otimes |1_{k\lambda}\rangle = |A^+\rangle \otimes |\psi_{\text{bound}}\rangle \otimes |1_{k\lambda}\rangle \equiv |A^+; \psi_{\text{bound}}\rangle |1_{k\lambda}\rangle$$

$$\equiv |A\rangle |1_{k\lambda}\rangle$$

$$|\phi_f\rangle \equiv |A^+\rangle \otimes |\psi_{\text{free}}\rangle \otimes |0_{k\lambda}\rangle \equiv |A^+; \psi_{\text{free}}\rangle |0_{k\lambda}\rangle$$

$|\psi_{\text{bound}}\rangle$ - initial state of electron in atom i.e. $|nlm\rangle$

$|\psi_{\text{free}}\rangle = |p_z^2/2m_e\rangle$ - final state of free electron

$$V_{fi} = (-e/m) \sqrt{\frac{\hbar}{2\epsilon_0 V \omega_k}} \langle \phi_f | \hat{a}_{k\lambda} e^{i\vec{k}\cdot\vec{r}} \hat{\epsilon}_{k\lambda} \cdot \vec{p} | \phi_i \rangle e^{-i\omega_k t} \quad (3)$$

↳ surviving term from $\sum_{k'\lambda'}$

$$C_{k\lambda} \equiv \langle \phi_f | \hat{a}_{k\lambda} e^{i\vec{k}\cdot\vec{r}} \hat{\epsilon}_{k\lambda} \cdot \vec{p} | \phi_i \rangle$$

$$= \langle A^+; \psi_{\text{free}} | \langle 0_{k\lambda} | \hat{a}_{k\lambda} e^{i\vec{k}\cdot\vec{r}} \hat{\epsilon}_{k\lambda} \cdot \vec{p} | A \rangle | 1_{k\lambda} \rangle$$

$$= \hat{\epsilon}_{k\lambda} \cdot \langle A^+; \psi_{\text{free}} | \langle 0_{k\lambda} | \hat{a}_{k\lambda} e^{i\vec{k}\cdot\vec{r}} \vec{p} | A \rangle | 1_{k\lambda} \rangle$$

~~Need $\langle \psi_{\text{free}} | \vec{p} | \psi_{\text{free}} \rangle$~~

$$= \hat{\epsilon}_{k\lambda} \cdot \underbrace{\langle A^+; \psi_{\text{free}} | e^{i\vec{k}\cdot\vec{r}} \vec{p} | A \rangle}_{\equiv \langle \text{atom} \rangle} \underbrace{\langle 0_{k\lambda} | \hat{a}_{k\lambda} | 1_{k\lambda} \rangle}_{\equiv \langle \text{rad} \rangle = \langle 0_{k\lambda} | 0_{k\lambda} \rangle = 1}$$

Need $e^{i\vec{k}\cdot\vec{r}} \vec{p} | \psi_{\text{free}} \rangle$ $\langle \psi_{\text{free}} | \hat{\epsilon}_{k\lambda} \cdot e^{i\vec{k}\cdot\vec{r}} \vec{p} = ?$

Drop \vec{k} subscript:

$$\begin{aligned} \vec{E} \cdot [\hat{\vec{p}}, e^{i\vec{k} \cdot \vec{r}}] \psi(\vec{r}) &= \vec{E} \cdot [\hat{\vec{p}} (e^{i\vec{k} \cdot \vec{r}} \psi) - e^{i\vec{k} \cdot \vec{r}} \hat{\vec{p}} \psi] \\ &= \vec{E} \cdot [-i\hbar \nabla (e^{i\vec{k} \cdot \vec{r}} \psi) - (-i\hbar) e^{i\vec{k} \cdot \vec{r}} \nabla \psi] \quad \text{using} \\ &= -i\hbar \vec{E} \cdot (i\vec{k} e^{i\vec{k} \cdot \vec{r}} \psi + e^{i\vec{k} \cdot \vec{r}} \nabla \psi - e^{i\vec{k} \cdot \vec{r}} \nabla \psi) \\ &= \hbar \vec{E} \cdot \vec{k} e^{i\vec{k} \cdot \vec{r}} \psi(\vec{r}) \\ &\quad \text{in Coulomb Gauge } \left(\vec{\nabla} \cdot \vec{A} = 0 \Rightarrow i\vec{k} \cdot \vec{A} = 0 \right. \\ &\quad \left. \Rightarrow \vec{k} \cdot \vec{E} = 0 \right) \\ &= 0 \end{aligned}$$

$$\Rightarrow \langle \psi_{\text{free}} | \vec{E} \cdot e^{i\vec{k} \cdot \vec{r}} \hat{\vec{p}} | \psi_{\text{free}} \rangle \neq 0$$

$$\begin{aligned} \langle \psi_{\text{free}} | \vec{E}_{\vec{k}\lambda} \cdot e^{i\vec{k} \cdot \vec{r}} \hat{\vec{p}} | \psi_{\text{free}} \rangle &= \langle \psi_{\text{free}} | \vec{E}_{\vec{k}\lambda} \hat{\vec{p}} e^{i\vec{k} \cdot \vec{r}} | \psi_{\text{free}} \rangle \\ &= \vec{E}_{\vec{k}\lambda} \cdot \langle \psi_{\text{free}} | \hat{\vec{p}} e^{i\vec{k} \cdot \vec{r}} | \psi_{\text{free}} \rangle = \vec{E}_{\vec{k}\lambda} \cdot \langle \frac{p e^2}{2m} | \hat{\vec{p}} e^{i\vec{k} \cdot \vec{r}} | \psi_{\text{free}} \rangle \\ &= \vec{E}_{\vec{k}\lambda} \cdot \vec{p}_e \langle \frac{p e^2}{2m} | e^{i\vec{k} \cdot \vec{r}} | \psi_{\text{free}} \rangle = \vec{E}_{\vec{k}\lambda} \cdot \vec{p}_e \langle \psi_{\text{free}} | e^{i\vec{k} \cdot \vec{r}} | \psi_{\text{free}} \rangle \end{aligned}$$

acts on bound electron in $|A\rangle$ i.e. on $|n, l, m\rangle$

$$\rightarrow C_{\vec{k}\lambda} = \vec{E}_{\vec{k}\lambda} \cdot \langle \text{atom} \rangle \cdot 1$$

$$\begin{aligned} \vec{E}_{\vec{k}\lambda} \cdot \langle A^+; \psi_{\text{free}} | e^{i\vec{k} \cdot \vec{r}} | A \rangle &= \vec{E}_{\vec{k}\lambda} \cdot \vec{p}_e \langle A^+; \psi_{\text{free}} | e^{i\vec{k} \cdot \vec{r}} | A \rangle \\ &= \vec{E}_{\vec{k}\lambda} \cdot \vec{p}_e \langle A^+; \psi_{\text{free}} | e^{i\vec{k} \cdot \vec{r}} | A^+; \psi_{\text{bound}} \rangle \\ &= \vec{E}_{\vec{k}\lambda} \cdot \vec{p}_e \langle \psi_{\text{free}} | e^{i\vec{k} \cdot \vec{r}} | \psi_{\text{bound}} \rangle \underbrace{\langle A^+ | A^+ \rangle}_{=1} \end{aligned}$$

$$\begin{aligned}
 \langle \psi_{free} | e^{i\vec{k}\cdot\vec{r}} | \psi_{bound} \rangle &= \langle \psi_{free} | e^{i\vec{k}\cdot\vec{r}} | nlm \rangle \\
 &= \int d^3r \langle \psi_{free} | \vec{r} X \vec{r} | e^{i\vec{k}\cdot\vec{r}} | nlm \rangle \\
 &= \int d^3r \psi_{free}^*(\vec{r}) \psi_{nlm}(\vec{r}) e^{i\vec{k}\cdot\vec{r}} \\
 \Rightarrow c_{k\lambda} &= (\vec{\epsilon}_{k\lambda} \cdot \vec{p}_e) \int d^3r \psi_{free}^*(\vec{r}) \psi_{nlm}(\vec{r}) e^{i\vec{k}\cdot\vec{r}}
 \end{aligned}$$

~~initial~~

→ (3):

$$V_{fi} = \frac{e}{m} \left(\frac{\hbar}{2\epsilon_0 V \omega_k} \right)^{1/2} \int d^3r \underbrace{\psi_{free}^*(\vec{r})}_{\text{final}} \underbrace{\psi_{nlm}(\vec{r})}_{\text{initial}} e^{i\vec{k}\cdot\vec{r}} (\vec{\epsilon}_{k\lambda} \cdot \vec{p}_e) e^{-i\omega_k t}$$

Let $\psi_{nlm}(\vec{r}) = \psi_{100}(\vec{r})$ i.e. atom is initially in ground state hydrogenic

initial: $\psi_{100}(\vec{r}) = \frac{1}{\sqrt{\pi}} \left(\frac{Z}{a_0} \right)^{3/2} e^{-r/a_0}$ // initial (4)

final: $\psi_{free}(\vec{r}) = \frac{e^{i\vec{p}_e \cdot \vec{r} / \hbar}}{\sqrt{V}}$ // final free electron
 ↳ assume $E_{photon} \approx \hbar\omega \gg E_{ionization}$

$$|V_{fi}|^2 = \left(\frac{e}{m} \right)^2 \frac{\hbar}{2\epsilon_0 V \omega_k} \frac{Z^3}{\pi a_0^3} (\vec{\epsilon}_{k\lambda} \cdot \vec{p}_e)^2 \left| \int d^3r e^{-Zr/a_0} e^{i(\vec{k} - \vec{p}_e/\hbar) \cdot \vec{r}} \right|^2 \quad (5)$$

$$\Rightarrow d\Gamma_{fi} = \frac{2\pi}{\hbar} \left(\frac{e}{m} \right)^2 \frac{\hbar Z^3}{2\epsilon_0 \pi a_0^3} \frac{m_e p_e dE_e}{(2\pi\hbar)^3 V \omega} (\vec{\epsilon}_{k\lambda} \cdot \vec{p}_e)^2 \left| \int d^3r e^{-Zr/a_0} e^{i(\vec{k} - \vec{p}_e/\hbar) \cdot \vec{r}} \right|^2 \cdot d\Omega_e \delta(E_A + E_e - (E_A + \hbar\omega)) \quad (5')$$

$$\begin{aligned}
 d\Gamma'_{fi} &= \frac{2\pi}{\hbar} d\Omega_e \int dE_e |V_{fi}|^2 P(E_e) \delta(E_e + E_{A+} - E_A - \hbar\omega) \\
 &= \frac{2\pi}{\hbar} d\Omega_e \left\{ |V_{fi}|^2 P(E_e) \right\}_{E_e = P_e^2/2m_e \stackrel{!}{=} -(E_{A+} - E_A - \hbar\omega)} \\
 &= \frac{2\pi}{\hbar} d\Omega_e \frac{\hbar(e/m_e)^2}{2\epsilon_0 \pi a_0^3} \frac{m_e Z^3}{(2\pi\hbar)^3 \omega V} \left\{ P_e (\vec{E} \cdot \vec{p}_e)^2 |\dot{I}|^2 \right\}_{P_e \stackrel{!}{=} \sqrt{2m_e(\hbar\omega - E_{ioniz})}} \\
 &= d\Omega_e \frac{e^2/m_e}{\epsilon_0 a_0^3 V} \frac{Z^3}{(2\pi\hbar)^3 \omega} \left\{ P_e (\vec{E} \cdot \vec{p}_e)^2 |\dot{I}|^2 \right\} \quad (6) \\
 \text{where } \dot{I} &= \int d^3r e^{-Zr/a_0} i(\vec{k} - \vec{p}_e/\hbar) \cdot \vec{r}
 \end{aligned}$$

The dependence on the arbitrary "interaction volume" V is reconciled by resorting to the differential cross-section for the PE effect:

Total cross-section:

$$\sigma = \int \left(\frac{d\sigma}{d\Omega_e} \right) d\Omega_e$$

Det

$$\left(\frac{d\sigma}{d\Omega_e} \right)_{\hbar\omega \gg E_{ioniz}} \approx \frac{V}{c} \frac{d\Gamma'_{fi}}{d\Omega_e}$$

$$= \frac{e^2/m_e}{c\epsilon_0 a_0^3} \frac{Z^3}{(2\pi\hbar)^3 \omega} \left\{ P_e (\vec{E} \cdot \vec{p}_e)^2 |\dot{I}|^2 \right\}_{P_e \stackrel{!}{=} \sqrt{2m_e(\hbar\omega - E_{ioniz})}} \quad (7)$$

$$= \dots = 32 Z^5 a_0^2 \left\{ \left(\frac{P_e c}{\hbar\omega} \right) \left(\frac{\vec{E} \cdot \vec{p}_e}{m_e c} \right)^2 \frac{1}{(Z^2 + a_0^2 q^2)^4} \right\}_{P_e \stackrel{!}{=} \sqrt{\dots}}$$

or

For hydrogenic atoms:

$$\left(\frac{d\sigma}{d\Omega_e} \right) = 32 Z^5 a_0^2 \left\{ \left(\frac{P_e c}{\hbar\omega} \right) \left(\frac{\vec{E} \cdot \vec{p}_e}{m_e c} \right)^2 \frac{1}{(Z^2 + a_0^2 q^2)^4} \right\}_{P_e \stackrel{!}{=} \sqrt{\dots}} \quad (8)$$

For (8), I used the following:

$$a_0 = \frac{\hbar}{Z\alpha m_e c}, \quad \alpha = \frac{e^2}{4\pi\epsilon_0 \hbar c}, \quad \dot{I} = \frac{8\pi a_0^3 Z}{(Z^2 + a_0^2 q^2)^2} \quad (q^2 \equiv |\vec{k} - \vec{p}_e/\hbar|^2)$$

(see next page)

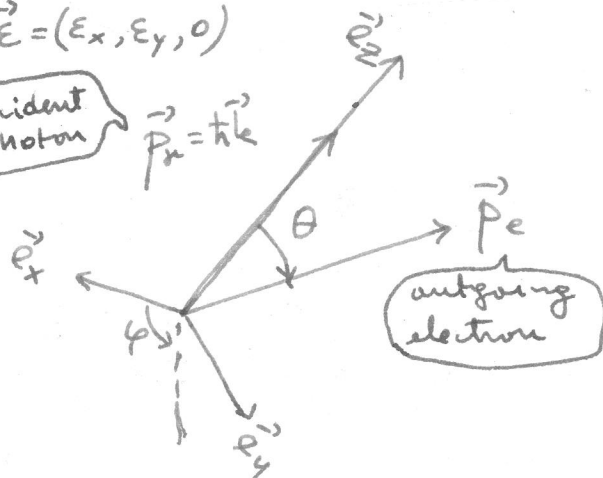
Evaluate " $\vec{E}_{\lambda} \cdot \vec{p}_e$ " from (5):

(Drop " λ " subscript $\vec{E}_{\lambda} \equiv \vec{E}$). Let photon direction be \vec{e}_z i.e. $\vec{k} = k \vec{e}_z$. Then polarization vector \vec{E} is in xy plane: $\vec{E} = (E_x, E_y, 0)$

$$\vec{p}_e = p_e (\sin \theta \cos \phi, \sin \theta \sin \phi, \cos \theta)$$

$$\vec{E} \cdot \vec{p}_e = E_x \sin \theta \cos \phi + E_y \sin \theta \sin \phi$$

incident photon $\vec{p}_\gamma = \hbar \vec{k}$



outgoing electron \vec{p}_e

* Linearly polarized incident photon:

$$\vec{E} = \vec{e}_x \Rightarrow \vec{E} \cdot \vec{p}_e = (\sin \theta \cos \phi) p_e$$

$$|V_{fi}|^2 = \frac{(e/m)^2 \hbar}{\pi a_0^3 V 2 \epsilon_0 \omega} (\sin^2 \theta \cos^2 \phi) \frac{8 \pi a_0^3 p_e^2 Z}{\left[1 + \left(a_0 \frac{p_e}{\hbar} \right)^2 \left(1 - \frac{v_e}{c} \cos \theta \right) \right]^2} \quad (p_e = \sqrt{2m_e E_e})$$

$$\Rightarrow \left[\frac{d\Gamma}{d\Omega_e} = \frac{V}{c} \frac{d\Gamma'_{fi}}{d\Omega_e} = \frac{2\pi}{\hbar} \frac{V}{c} \left\{ |V_{fi}|^2 \rho(p_e^2/2m) \right\} \right]_{p_e = \sqrt{2m(\hbar\omega - E_{ion})}}$$

* Circularly polarized incident photon: (i.e. unpolarized)

$$\vec{E} = \vec{E}_1 = 1/\sqrt{2} (\vec{e}_x + i \vec{e}_y) \Rightarrow \vec{E} \cdot \vec{p}_e = 1/\sqrt{2} (\sin \theta \cos \phi + i \sin \theta \sin \phi) p_e$$

$$= \frac{\sin \theta}{\sqrt{2}} (\cos \phi + i \sin \phi) p_e$$

$$= \frac{p_e}{\sqrt{2}} \sin \theta e^{i\phi}$$

$$|V_{fi}|^2 = \frac{(e/m)^2 \hbar}{\pi a_0^3 V \epsilon_0 \omega} \frac{p_e^2 \sin^2 \theta}{4} |I|^2$$

$$\Rightarrow \left[\frac{d\Gamma}{d\Omega_e} = \frac{V}{c} \frac{d\Gamma'_{fi}}{d\Omega_e} \approx \dots Z^7 2\sqrt{2} \alpha^8 \frac{a_0^2}{Z^2} \left(\frac{Mc^2}{E_e} \right)^{7/2} \frac{\sin^2 \theta}{\left(1 - \frac{v_e}{c} \cos \theta \right)^4} \right]$$

Premise:

$$\hbar\omega \gg E_{ioniz} (= E_{At} - E_A) \Leftrightarrow E_e \gg \alpha^2 \frac{Mc^2}{2}$$